



Mixed Finite Elements for Variational Surface Modeling

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Motivation

Produce high-quality surfaces

Via energy minimization

Or solving Partial Differential Equations

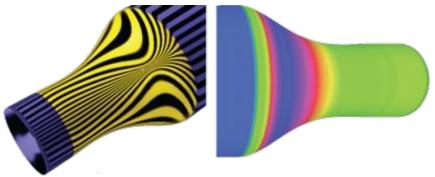
$$\langle f,g \rangle = \int fg$$

Laplacian energy

$$E_B = \frac{1}{2} \langle \Delta \mathbf{u}, \Delta \mathbf{u} \rangle_{\Omega_0} \to \min$$

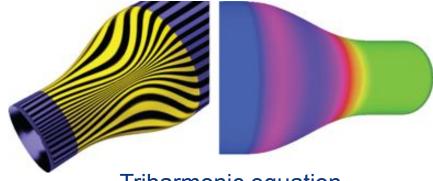
Laplacian gradient energy

$$E_B = \frac{1}{2} \langle \Delta \mathbf{u}, \Delta \mathbf{u} \rangle_{\Omega_0} \to \min \quad E_T = \frac{1}{2} \langle \nabla \Delta \mathbf{u}, \nabla \Delta \mathbf{u} \rangle_{\Omega_0} \to \min$$



Biharmonic equation

$$\Delta^2 \mathbf{u} = 0$$

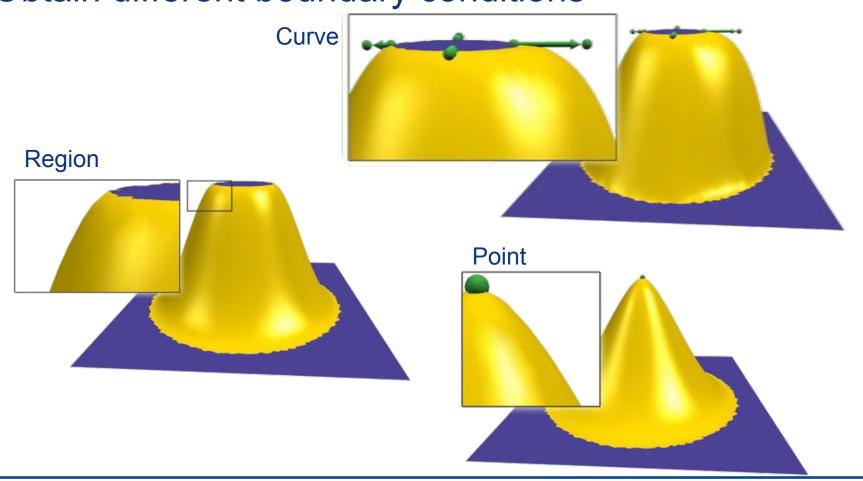


Triharmonic equation

$$\Delta^3 \mathbf{u} = 0$$

Motivation

Obtain different boundary conditions



Motivation

Higher-order equations on mesh (i.e. piecewise linear elements)

Dealing with higher-order derivatives not straightforward

Previous work

Simple domains, analytic boundaries

[Bloor and Wilson 1990]

Model shaped minimization of curvature variation energy

[Moreton and Séquin 1992]

Interpolate curve networks, local quadratic fits and finite differences

[Welch and Witkin 1994]

Uniform-weight discrete Laplacian

[Taubin 1995]

Cotangent-weight discrete Laplacian

[Pinkall and Polthier 1993],

[Wardetzky et al. 2007],

[Reuter et al. 2009]

Wilmore flow, using FEM with aux variables

Position and co-normal specification on boundary

[Clarenz et al. 2004]

Linear systems for k-harmonic equations

Uses discretized Laplacian operator

[Botsch and Kobbelt 04]

Previous work

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Standard Finite Element Method

Requires at least C¹ elements for fourth order

Can't use standard triangle meshes

High order surfaces exist, (e.g. Argyris triangle)

- Require many extra degrees of freedom
- Not popular due to complexity

Low order, C⁰, workarounds

E.g. mixed elements

Discrete Geometric Discretization

Derive mesh analog of geometric quantity

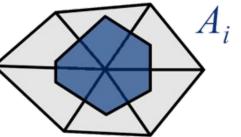
E.g. Laplace-Beltrami operator integrated over vertex area

Re-expressed using only first-order derivatives

Use average value as energy of vertex area

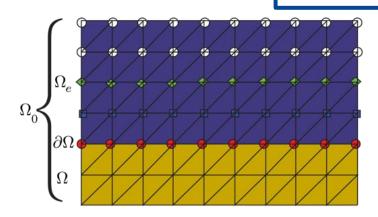
Used often in geometric modeling

No obvious way to connect to continuous case



Introduce additional variable to convert high order problem to low order

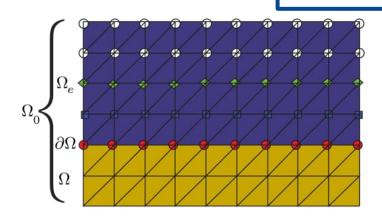
$$\frac{1}{2}\langle \Delta \mathbf{u}, \Delta \mathbf{u} \rangle_{\Omega_0} \to \min$$
 \to



$$\frac{1}{2}\langle \Delta \mathbf{u}, \Delta \mathbf{u} \rangle_{\Omega_0} \to \min \quad \Rightarrow \quad \frac{1}{2}\langle \mathbf{v}, \mathbf{v} \rangle_{\Omega_0} \to \min, \text{ s.t. } \Delta \mathbf{u} = \mathbf{v}$$

 $\Delta^2 \mathbf{u} = 0$

Introduce additional variable to convert high order problem to low order



$$\frac{1}{2}\langle \Delta \mathbf{u}, \Delta \mathbf{u} \rangle_{\Omega_0} \to \min \quad \rightarrow \quad \frac{1}{2}\langle \mathbf{v}, \mathbf{v} \rangle_{\Omega_0} \to \min, \text{ s.t. } \Delta \mathbf{u} = \mathbf{v}$$

Use Langrange multipliers to enforce constraint

$$L_B = \frac{1}{2} \langle \mathbf{v}, \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \Delta \mathbf{u} - \mathbf{v} \rangle_{\Omega_0} = \frac{1}{2} \langle \mathbf{v}, \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \frac{\partial \mathbf{u}}{\partial n} \rangle_{\partial \Omega_0} - \langle \nabla \lambda, \nabla \mathbf{u} \rangle_{\Omega_0}$$

Constraint structure also makes certain boundary types easier

$$\Delta^2 \mathbf{u} = 0$$

Our original higher order problem

$$\Delta^2 \mathbf{u} = 0$$

Introduce an additional variable

$$\Delta \mathbf{u} = \mathbf{v}$$

$$\Delta \mathbf{v} = 0$$

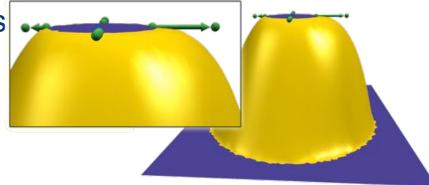
 ∂n

Two second order problems

Can use just linear elements

Curve

- lacksquare Fixed boundary curve $\,\partial {f u}$
- Specified tangents:



Discretize with piecewise linear approximations for variables

$$\sum_{j \in I} (\mathbf{v}_j - \lambda_j) \langle \phi_j, \phi_i \rangle_{\Omega_0} = 0$$

$$-\sum_{j\in I} \lambda_j \langle \nabla \phi_j, \nabla \phi_j \rangle_{\Omega_0} = 0$$

$$\sum_{j \in I} \mathbf{v}_j \langle \phi_j, \phi_i \rangle_{\Omega_0} + \sum_{j \in I} \frac{\partial \mathbf{u}_j}{\partial n} \langle \phi_j, \phi_i \rangle_{\partial \Omega_0} - \sum_{j \in I} \mathbf{u}_j \langle \nabla \phi_j, \nabla \phi_i \rangle_{\Omega_0} = 0$$

Discrete Laplacian
$$L_{ij} = -\langle \nabla \phi_i, \nabla \phi_j \rangle$$

Mass matrix
$$M_{ij}^{
m full} = -\langle \phi_i, \phi_j
angle$$

$$\Delta^2 \mathbf{u} = 0$$

Discretize with piecewise linear approximations for variables

$$\sum_{j \in I} (\mathbf{v}_j - \lambda_j) \langle \phi_j, \phi_i \rangle_{\Omega_0} = 0$$

$$-\sum_{j\in I} \lambda_j \langle \nabla \phi_j, \nabla \phi_j \rangle_{\Omega_0} = 0$$

$$\sum_{j \in I} \mathbf{v}_j \langle \phi_j, \phi_i \rangle_{\Omega_0} + \sum_{j \in I} \frac{\partial \mathbf{u}_j}{\partial n} \langle \phi_j, \phi_i \rangle_{\partial \Omega_0} - \sum_{j \in I} \mathbf{u}_j \langle \nabla \phi_j, \nabla \phi_i \rangle_{\Omega_0} = 0$$

Discrete Laplacian
$$L_{ij} = -\langle \nabla \phi_i, \nabla \phi_j \rangle$$

Mass matrix

$$M_{ij}^{\mathrm{full}} pprox M_{ij}^{\mathrm{d}}$$

Matrix form, curve boundary conditions

$$\begin{bmatrix} -M^{\Omega} & L_{\bar{\Omega},\Omega} \\ L_{\Omega,\bar{\Omega}} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{u}_{\Omega} \end{bmatrix} = \begin{bmatrix} -L_{\bar{\Omega},0}^{\Omega} \mathbf{b}_0 - N_{\bar{\Omega},0}^{\partial \Omega} \mathbf{b}_1 \\ 0 \end{bmatrix}$$

$$L_{ij} = -\langle \nabla \phi_i, \nabla \phi_j \rangle$$

$$M_{ij}^{\mathrm{d}} \approx -\langle \phi_j, \phi_i \rangle$$

Discrete Laplacian
$$L_{ij}=-\langle
abla \phi_i,
abla \phi_j
angle$$
 Mass matrix $M_{ij}^{\mathrm{d}} pprox -\langle \phi_j, \phi_i
angle$ Neumann matrix $N_{ij}^{\partial\Omega}=\langle \phi_j, \phi_i
angle_{\partial\Omega}$

Where
$$\mathbf{u}=\mathbf{b}_0$$
 and $\dfrac{\partial \mathbf{u}}{\partial n}=\mathbf{b}_n$

Diagonalized, lumped mass matrices eliminate auxiliary variable

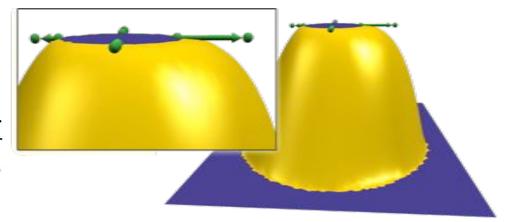
$$L_{\Omega,\bar{\Omega}}(M^{\mathrm{d}})^{-1}L_{\bar{\Omega},\Omega}\mathbf{u}_{\Omega} = -L_{\Omega,\bar{\Omega}}(M^{\mathrm{d}})^{-1}(-L_{\bar{\Omega},0}^{\Omega}\mathbf{b}_{0} - N_{\bar{\Omega},0}^{\partial\Omega}\mathbf{b}_{1})$$

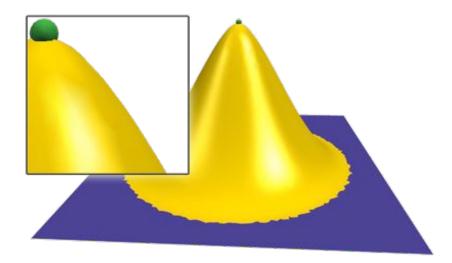
Curve

- Fixed boundary curve
- lacksquare Specified tangents: $\partial \mathbf{u}$ ∂n

Point

Single fixed points on surface

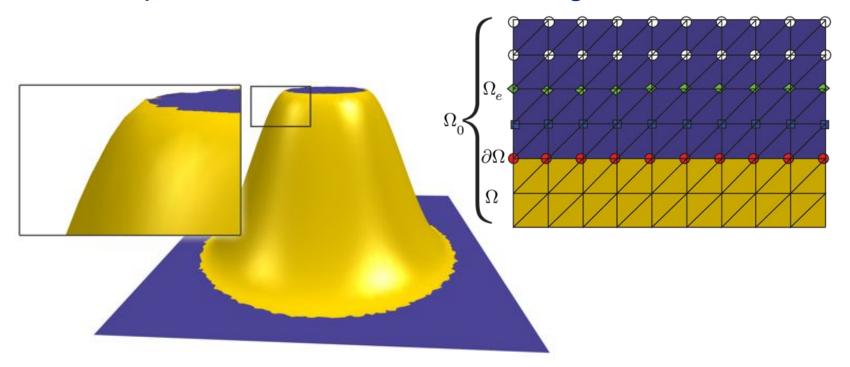




$\Delta^2 \mathbf{u} = 0$

Region

Fixed part of mesh outside solved region



Use Lagrangian to enforce region condition

$$L_B = \frac{1}{2} \langle \mathbf{v}, \mathbf{v} \rangle_{\Omega_0} - \langle \lambda, \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \frac{\partial \mathbf{u}}{\partial n} \rangle_{\partial \Omega} - \langle \nabla \lambda, \nabla \mathbf{u} \rangle_{\Omega_0} + \langle \mu, \mathbf{u}_f - \mathbf{u} \rangle_{\Omega_f}$$

Discretize with piecewise linear approximations for variables

$$\begin{bmatrix} -M^d & L_{\bar{\Omega},\Omega} \\ L_{\Omega,\bar{\Omega}} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{u}_{\Omega} \end{bmatrix} = \begin{bmatrix} -L_{\bar{\Omega},0}\mathbf{u}_0^f - L_{\bar{\Omega},1}\mathbf{u}_1^f \\ 0 \end{bmatrix}$$
 Discrete Laplacian
$$L_{ij} = -\langle \nabla \phi_i, \nabla \phi_j \rangle$$
 Mass matrix

May also eliminate aux. variable

Mass matrix
$$M_{ij}^{
m d} pprox -\langle \phi_j, \phi_i
angle$$

$$L_{\Omega,\bar{\Omega}}(M^{\mathrm{d}})^{-1}L_{\bar{\Omega},\Omega}\mathbf{u}_{\Omega} = -L_{\Omega,\bar{\Omega}}(M^{\mathrm{d}})^{-1}L_{\bar{\Omega},01}\mathbf{u}^{f}$$

$$\Delta^2 \mathbf{u} = 0$$

$$\begin{bmatrix} -\mathbf{M}^d & L_{\bar{\Omega},\Omega} \\ L_{\Omega,\bar{\Omega}} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{u}_{\Omega} \end{bmatrix} = \begin{bmatrix} -L_{\bar{\Omega},0}\mathbf{u}_0^f - L_{\bar{\Omega},1}\mathbf{u}_1^f \\ 0 \end{bmatrix} \begin{bmatrix} -\mathbf{M}^{\Omega} & L_{\bar{\Omega},\Omega} \\ L_{\Omega,\bar{\Omega}} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{u}_{\Omega} \end{bmatrix} = \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} = \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} = \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} = \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} = \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} = \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{v}_{\bar{\Omega}} \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega$$

$$\begin{bmatrix} -M^{\Omega} & L_{\bar{\Omega},\Omega} \\ L_{\Omega,\bar{\Omega}} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{u}_{\Omega} \end{bmatrix} = \begin{bmatrix} -L_{\bar{\Omega},0}^{\Omega} \mathbf{b}_0 - N_{\bar{\Omega},0}^{\partial \Omega} \mathbf{b}_1 \\ 0 \end{bmatrix}$$

Difference in right-hand side

Curve conditions don't require lumped mass matrix

But we use it in practice, for speed and numerical accuracy

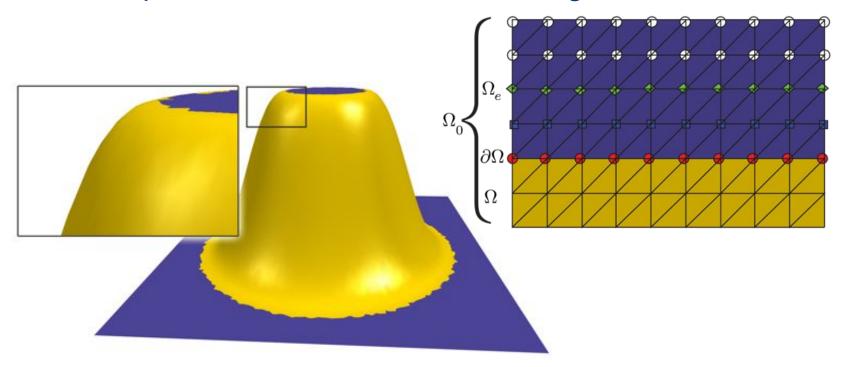
Equivalent to [Botsch and Kobbelt, 2004]

■ Specified tangents ≈ parameter for continuity control

$\Delta^3 \mathbf{u} = 0$

Region

Fixed part of mesh outside solved region



$$\Delta^3 \mathbf{u} = 0$$

Convert high order problem to low order problem

$$\frac{1}{2}\langle \nabla \Delta \mathbf{u}, \nabla \Delta \mathbf{u} \rangle_{\Omega_0} \to \min \quad \Rightarrow \quad \frac{1}{2}\langle \nabla \mathbf{v}, \nabla \mathbf{v} \rangle_{\Omega_0} \to \min, \text{ s.t. } \Delta \mathbf{u} = \mathbf{v}$$

Use Langrange multipliers to enforce constraint

$$L_T = \frac{1}{2} \langle \nabla \mathbf{v}, \nabla \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \Delta \mathbf{u} - \mathbf{v} \rangle_{\Omega_0} =$$

$$\frac{1}{2}\langle \nabla \mathbf{v}, \nabla \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \frac{\partial \mathbf{u}}{\partial n} \rangle_{\partial \Omega_0} - \langle \nabla \lambda, \nabla \mathbf{u} \rangle_{\Omega_0}$$

$$\Delta^3 \mathbf{u} = 0$$

Convert high order problem to low order problem

$$\frac{1}{2}\langle\nabla\Delta\mathbf{u},\nabla\Delta\mathbf{u}\rangle_{\Omega_0}\rightarrow\min\quad \boldsymbol{\rightarrow}\ \frac{1}{2}\langle\nabla\mathbf{v},\nabla\mathbf{v}\rangle_{\Omega_0}\rightarrow\min,\ \mathrm{s.t.}\ \Delta\mathbf{u}=\mathbf{v}$$

Use Langrange multipliers to enforce constraint

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$$\frac{1}{2}\langle \nabla \mathbf{v}, \nabla \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \frac{\partial \mathbf{u}}{\partial n} \rangle_{\partial\Omega_0} - \langle \nabla \lambda, \nabla \mathbf{u} \rangle_{\Omega_0}$$

Notice similarity to Lagrangian for biharmonic

$$L_B = \frac{1}{2} \langle \mathbf{v}, \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \Delta \mathbf{u} - \mathbf{v} \rangle_{\Omega_0} =$$

$$\frac{1}{2}\langle \mathbf{v}, \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \mathbf{v} \rangle_{\Omega_0} + \langle \lambda, \frac{\partial \mathbf{u}}{\partial n} \rangle_{\partial\Omega_0} - \langle \nabla \lambda, \nabla \mathbf{u} \rangle_{\Omega_0}$$

$$\Delta^3 \mathbf{u} = 0$$

Discretization, formulation works the same way

$$\begin{bmatrix} L_{\bar{\Omega}\bar{\Omega}} & 0 & -M^d_{\bar{\Omega}\bar{\Omega}} \\ 0 & 0 & L_{\Omega\bar{\Omega}} \\ -M^d_{\bar{\Omega}\bar{\Omega}} & L_{\bar{\Omega}\Omega} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{u}_{\Omega} \\ \lambda_{\bar{\Omega}} \end{bmatrix} = \begin{bmatrix} -L_{\bar{\Omega},1}\mathbf{v}_1 \\ 0 \\ -L_{\bar{\Omega},0}\mathbf{u}_0^f - L_{\bar{\Omega},1}\mathbf{u}_1^f \end{bmatrix}$$

where
$$\mathbf{v}_1 = -L_{1,0}\mathbf{u}_0^f - L_{1,1}\mathbf{u}_1^f - L_{1,2}\mathbf{u}_2^f$$

Eliminate auxiliary variables

Leaving system with only u

Discrete Laplacian $L_{ij} = -\langle \nabla \phi_i, \nabla \phi_j \rangle$

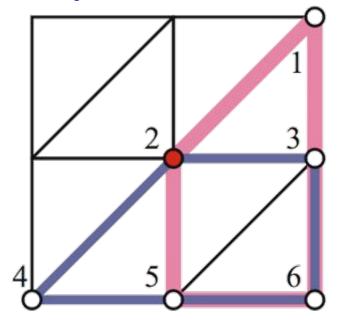
Mass matrix $M_{ij}^{\mathrm{full}} = -\langle \phi_i, \phi_j
angle$

 $\Delta^3 \mathbf{u} = 0$

Curve

- Fixed boundary curve
- Specified tangents and curvatures:

Leads to singular systems

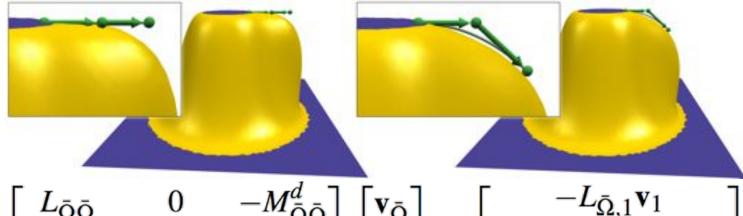


 $\Delta^3 \mathbf{u} = 0$

Curve → Region

- Fixed boundary curve and one ring into interior
- Specified curvatures: $\frac{\partial^2 \mathbf{u}}{\partial n^2}$

$$\mathbf{v} = \Delta \mathbf{u} = \frac{\partial^2 \mathbf{u}}{\partial n^2} + \frac{\partial^2 \mathbf{u}}{\partial t^2}$$



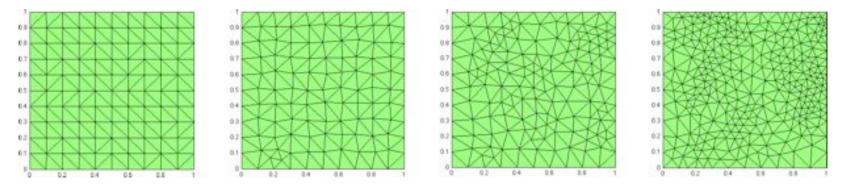
$$\begin{bmatrix} L_{\bar{\Omega}\bar{\Omega}} & 0 & -M^d_{\bar{\Omega}\bar{\Omega}} \\ 0 & 0 & L_{\Omega\bar{\Omega}} \\ -M^d_{\bar{\Omega}\bar{\Omega}} & L_{\bar{\Omega}\Omega} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{v}_{\bar{\Omega}} \\ \mathbf{u}_{\Omega} \\ \lambda_{\bar{\Omega}} \end{bmatrix} = \begin{bmatrix} -L_{\bar{\Omega},1}\mathbf{v}_1 \\ 0 \\ -L_{\bar{\Omega},0}\mathbf{u}_0^f - L_{\bar{\Omega},1}\mathbf{u}_1^f \end{bmatrix}$$

Experimental Results

Tested convergence of our systems

Randomly generated domains of varying irregularity

- One vertex placed randomly in each square of grid
- Parameter controlled variation from regular



Connected using Triangle Library

Control minimal interior angles

Experimental Results

Specify boundary conditions using analytic target functions: \mathbf{u}^t

Try to reproduce original function by solving system:

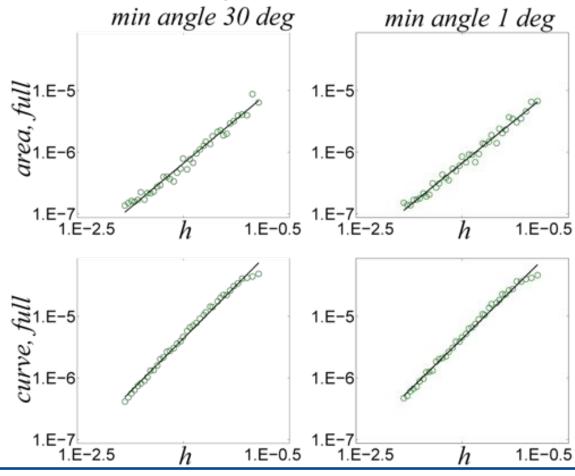
$$\Delta^2 \mathbf{u} = \Delta^2 \mathbf{u}^t$$

Measure error between analytic target and our mixed FEM approximation

$\Delta^2 \mathbf{u} = 0$

Experimental Results

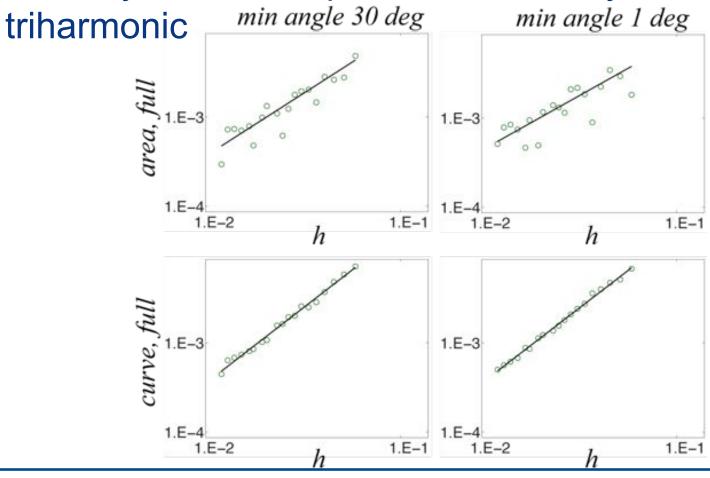
Nearly optimal convergence for biharmonic



$\Delta^3 \mathbf{u} = 0$

Experimental Results

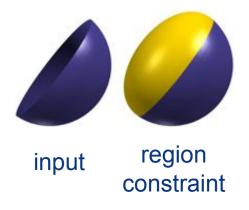
Boundary conditions perform differently for



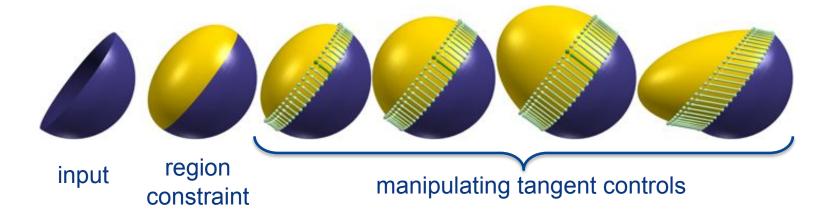
 $\Delta^2 \mathbf{u} = 0$



 $\Delta^2 \mathbf{u} = 0$



 $\Delta^2 \mathbf{u} = 0$



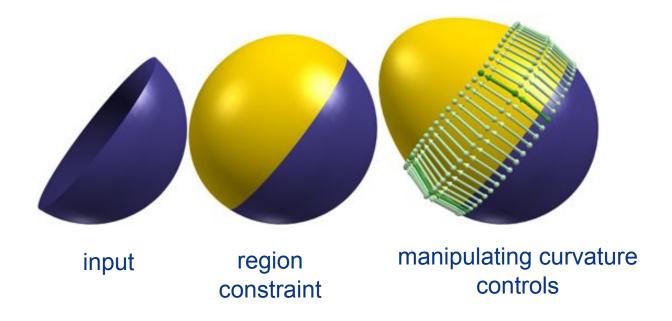
 $\Delta^2 \mathbf{u} = 0$



manipulating tangent controls

 $\Delta^3 \mathbf{u} = 0$

Filling in holes: Laplacian gradient energy

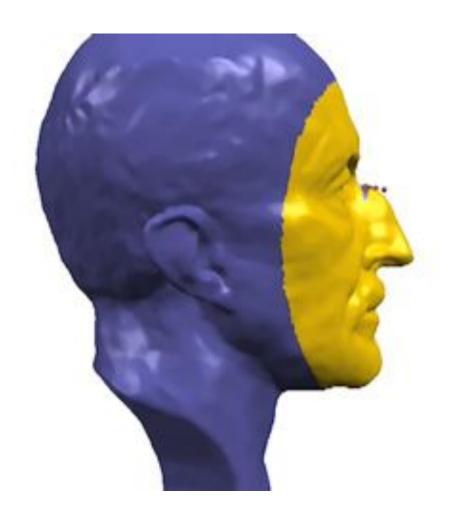


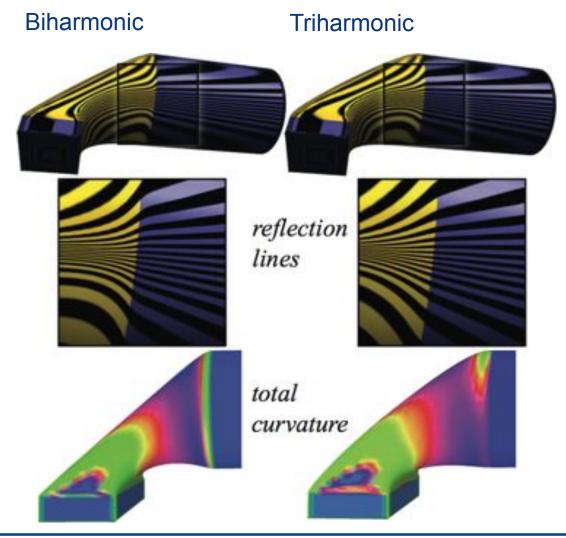
 $\Delta^2 \mathbf{u} = 0$

Specifying tangents in Laplacian energy around regions



$\Delta^3 \mathbf{u} = 0$





Summary

Technique for discretizing energies or PDEs

- Reduce to low order by introducing variables
- Use constraints to enforce region boundary conditions
- Lump mass matrix

Convergence for fourth- and sixth-order PDEs

Summary

Technique for discretizing energies or PDEs

- Reduce to low order by introducing variables
- Use constraints to enforce region boundary conditions
- Lump mass matrix

Convergence for fourth- and sixth-order PDEs

Future work

- Improve convergence of triharmonic solution
- Explore using non-flat metric

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